

Laser-based RF Precision Oscillators

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S. J. Augst, D. H. Spears, J. Twichell, C. Freed and T. Y. Fan MIT Lincoln Laboratory, Lexington, MA

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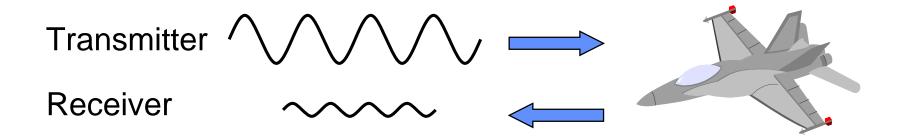


Outline

- Background Doppler radar
- CO₂ laser RF source
- Nd:YAG laser RF source
- Summary of proposed work



Doppler Shift



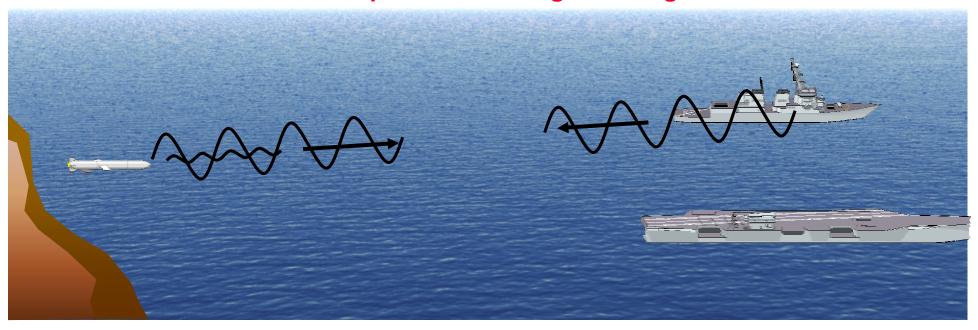
- Target moving towards you shifts return up in frequency, away => down
- Doppler shift is proportional to frequency = f * 2v/c
- Doppler shift often separates target from background clutter

	S-Band (3 GHz)	X-Band (10 GHz)
Crawl (1/4 m/s)	5 Hz	17 Hz
Walk (1 m/s)	20 Hz	67 Hz
Drive (10 m/s)	200 Hz	667 Hz
Fly (100 m/s)	2000 Hz	6.67 kHz
Mach 1 (340 m/s)	6.8 kHz	22.7 kHz
LEO (7000 m/s)	140 kHz	467 kHz



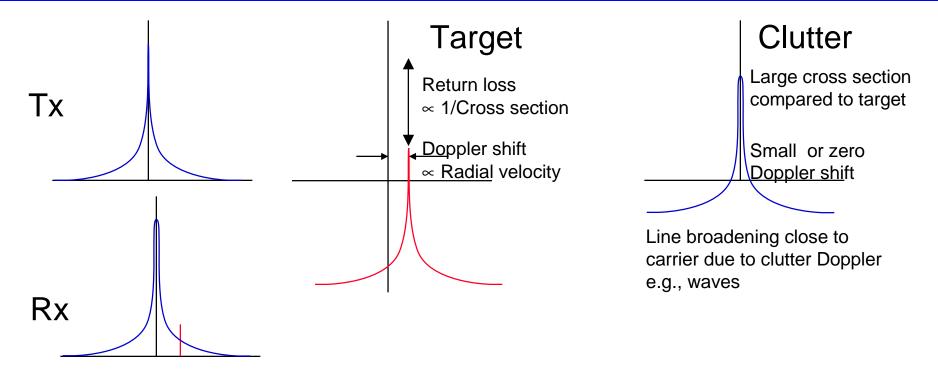
Targets Hide in Clutter

- Clutter (environmental returns) gives *huge* signal compared to target, can be > 10⁶ larger for some targets
- Target return is slightly Doppler shifted in frequency
- Clutter returns transmit signal and transmit noise
- Transmit noise in clutter return can hide target
- Better oscillator helps see small targets hiding in clutter





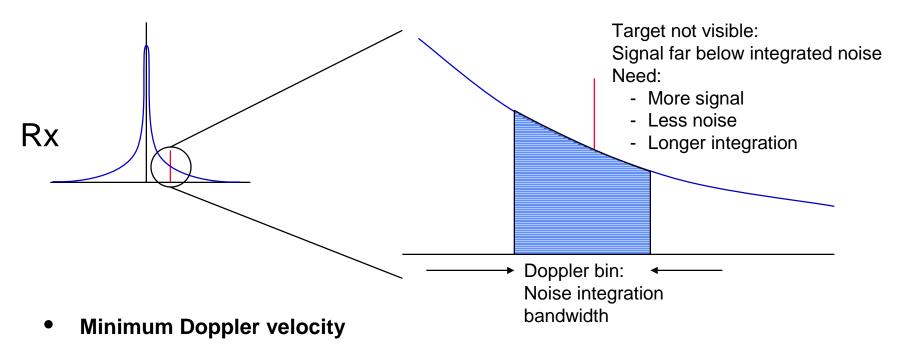
Oscillator noise in a Doppler radar system can obscure the target



- Oscillator noise is transmitted, blurring the return, and mixed into the return in the receiver.
- Noise raises the noise floor, particularly for small frequency shifts (small Doppler velocity)
- This limits the ability to see small targets with small velocity against clutter



Doppler Bin-width Sets Detection Limit



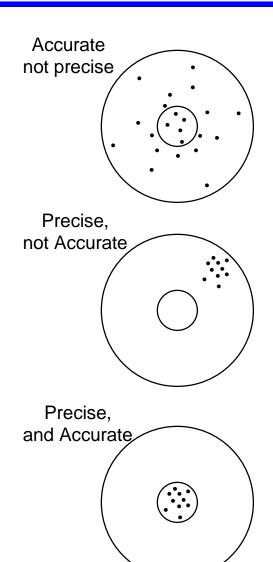
- Integration time
- Clutter doppler
- Phase noise
- Doppler bin width set to ~10% of Doppler frequency
- Sets clutter & phase noise bandwidth
- Clutter Improvement Factors (CIF) of 60 to 110 dB needed



State of the Art

- For times < 1 sec, Quartz is standard to beat
- Atomic clocks use "physics package" for long term accuracy
 - Terrible SNR
 - Shot noise limited with small number of photons
 - Accurate by virtue of fundamental physics
 - Long (>1000 sec) integration times
- "Flywheel" used for short term precision
 - Quartz
 - Ultra-stable cavity, CO₂ or Nd:YAG lasers

Path to long-term accuracy – locking laser to an atomic reference



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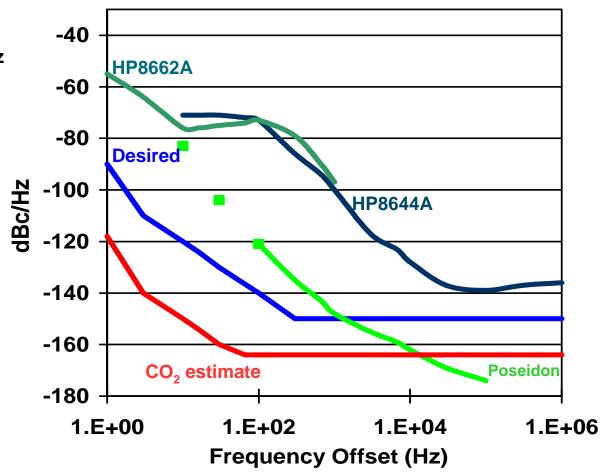


Desired Phase Noise Spectrum

Caveats

- All sources normalized to 2.7 GHz
 - 6 dB/octave ideal multiplier or divider
- "Desired" assumes
 - Ideal RF components
 - 12 bit digitizer limits SNR
- CO₂ phase noise estimate
 - Allen variance translation
 - 2 laser tests
 - 10 mW shot noise into ideal detector

Phase Noise & Doppler Shifts Normalized to 2.7 GHz





Schawlow-Townes linewidth gives the quantum noise limit for a laser cavity

$$\Delta v = \left(\frac{ahc^2}{4\pi}\right) \frac{T^2 v}{L^2 P_{out}}$$

T=round-trip cavity loss *v*=optical frequency *L*=cavity length P_{out} =output optical power *h*=Planck's constant *a*=inversion parameter (≈1) c=light speed

CO₂

T
$$pprox 3\%, \ Lpprox 0.5m,
ightarrow rac{\Delta \, \mathcal{V}}{\mathcal{V}} pprox 2x10^{-20} \ Ppprox 1W, \ \Delta \, \mathcal{V} pprox 5x10^{-7} \ Hz$$

Nd:YAG



Using a two-frequency laser as a quiet RF oscillator

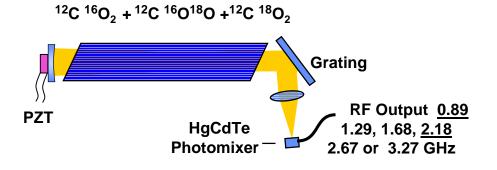
- A laser's optical frequency can be much quieter than stateof-the-art crystal oscillators.
- Optical frequencies (~100THz) cannot be detected by photodetectors, but a photodetector can see the beat frequency between two closely spaced optical frequencies.
- The optical frequency noise depends on laser cavity length instability which will map onto the RF beat frequency if the two frequencies are spatially overlapped in the same cavity:

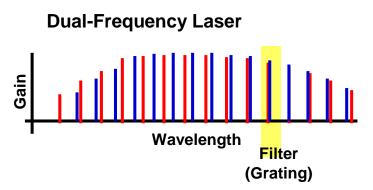
$$\frac{\Delta L}{L} = \frac{\Delta V_{opt}}{V_{opt}} \left(= \frac{\Delta f_{RF}}{f_{RF}} \right)$$

Single cavity eliminates many common-path errors



CO₂ -Laser-Based RF Frequency Reference

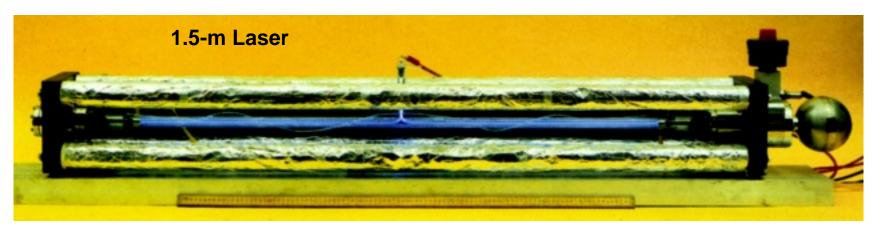




- RF beat note (f_{RF}) from dual-frequency, single-cavity laser
- Common optical cavity minimizes RF noise, drift, thermal effects, acoustic effects, etc.
- In 1979, $\Delta f_{RF}/f_{RF}$ was found to be less than that of hp 8672A synthesizer for times < 200 µs



Characteristics of Freed CO₂ Lasers



Features

Super invar spacer rods with thermal, vibration, acoustical & magnetic shielding Black Diabase granite end plates, Stabilized DC plasma discharge

Achieved $\Delta v/v < 2 \times 10^{-13}$

<u>Laser</u>	Gain Length (m)	Comments
1.5-m Grating Output Coupled	1.23	Long gain region, mixed isotopes ok
0.5-m Two-Mirror	0.23	High-reflectivity mirrors needed for mixed isotopes

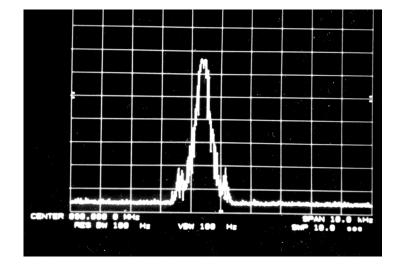


890 MHz Beat from Two-Frequency Mixed-Isotope CO₂ Laser

¹⁶O¹²C¹⁶O I P(12) and ¹⁶O¹²C¹⁸O I P(19) Lines

100 Hz Instrument Resolution



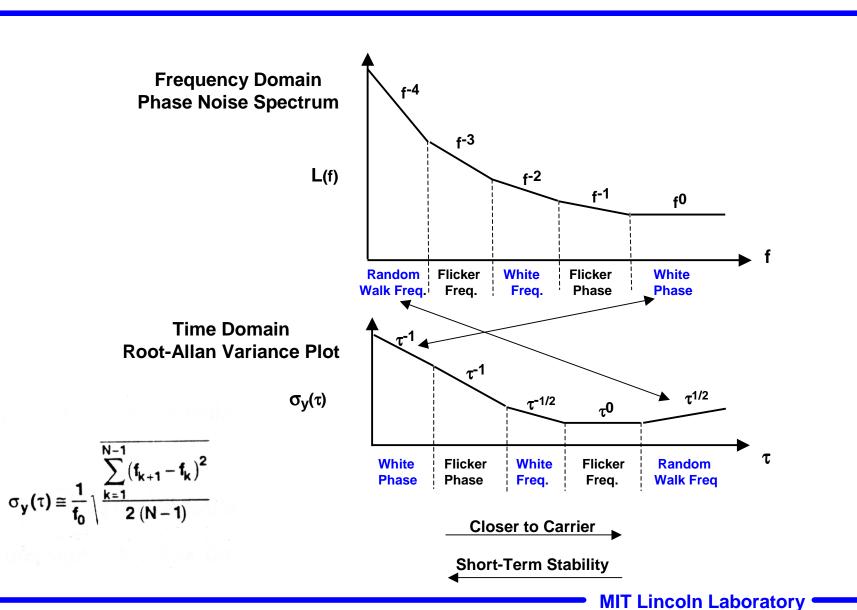


1 kHz/div

~0.3 mW of CO₂ Power on Photomixer -28 dBm from HgCdTe Photomixer +2 dBm with 30-dB amplifier



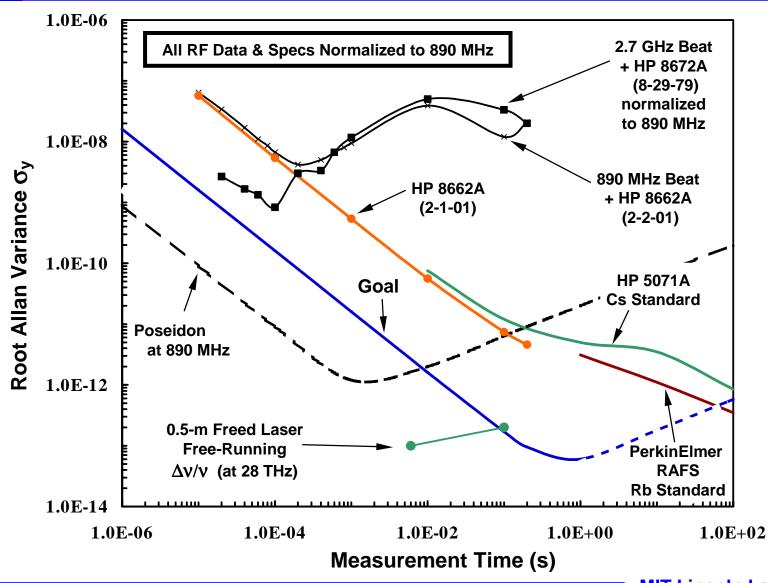
Phase Noise (f) and Allan Variance (τ)



14

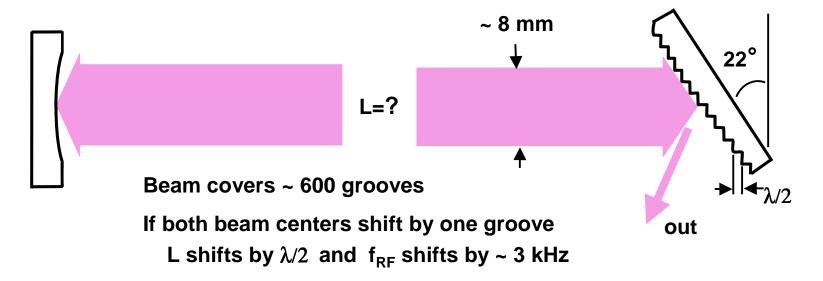


CO₂ Laser Root Allan Variance Data



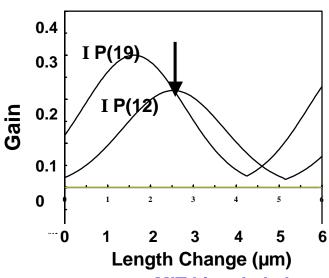


No Precise Cavity Length in Grating Laser



Gain slope at P(19) operating point pulling towards shorter cavity

Anomalously large variations in 890-MHz beat frequency are a result of 626 I P(12) and 628 I P(19) operating at slightly different cavity lengths (1 nm ~ 19 kHz)

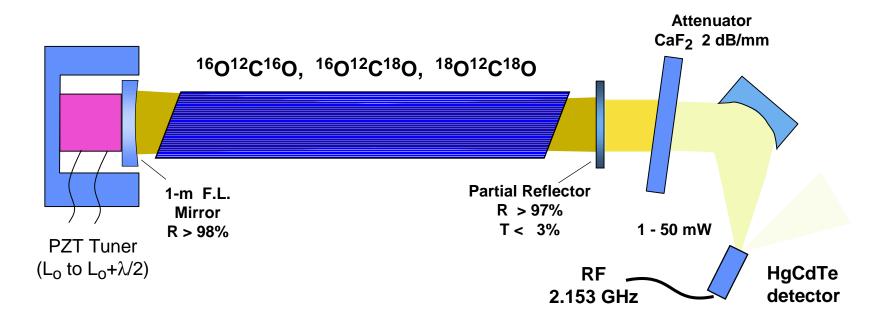


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Dual-Frequency, Two-Mirror Single-Cavity Laser

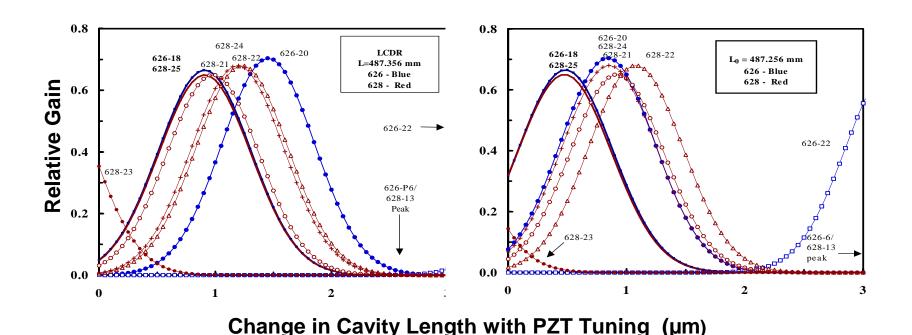
- Dual-frequency, single-cavity laser minimizes RF phase noise
- 0.5-m two-mirror Freed lasers have shown 1:10¹³ short-term stability
- Main Issue: How to get dual-frequency operation without a grating





Calc. Gain vs. Cavity Length for 0.5-m Laser

(Shown are only 626 I P18 and 628 I P25 transitions and those I P branch (10.6 µm) lines with higher gain)



At length ~ 100 µm shorter than LCDR, the 626 I P18 and 628 I P25 lines are well separated from other strong 10.6-µm lines



Criteria for Dual-Frequency Two-Mirror Mixed-Isotope CO₂ Laser

- Short cavity (<0.5 m, 300 MHz FSR) to allow ~10 lines
- Low pressure (<20 torr, <50 MHz FWHM) allows ~ 10 lines
- ¹6O/¹8O isotope ratio optimized to balance gains

The above leads to LOW GAIN

- High-reflectivity mirrors (>99% spherical and >98% output)
- Selective coated surface in cavity to reject 9.5-µm lines
- Precise cavity length for line-center double resonance
 - (e.g. 487.356 mm for 626 I P18 / 628 I P25 beat at 2.15 GHz)



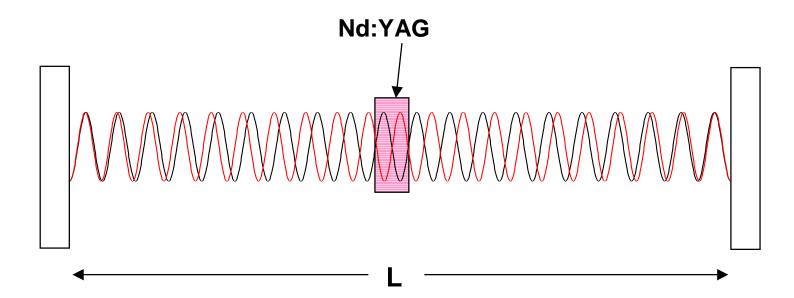
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Standing-wave cavity lasing in two modes

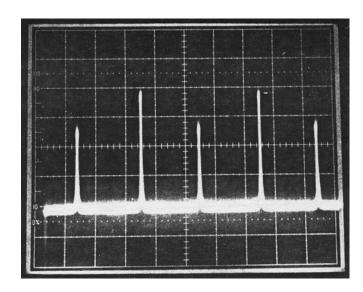
- Spatial hole-burning causes the laser cavity to oscillate in exactly two longitudinal modes.
- The beat frequency will be equal to the cavity freespectral-range (= c/2L).



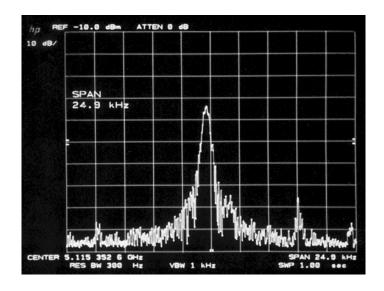


Two-mode Nd:YAG laser proof-of-principle experiment (1993)

- Two-mode lasing has been demonstrated using the spatialhole-burning technique.
- The initial demonstration was a table-top design, not intended to have excellent noise performance.





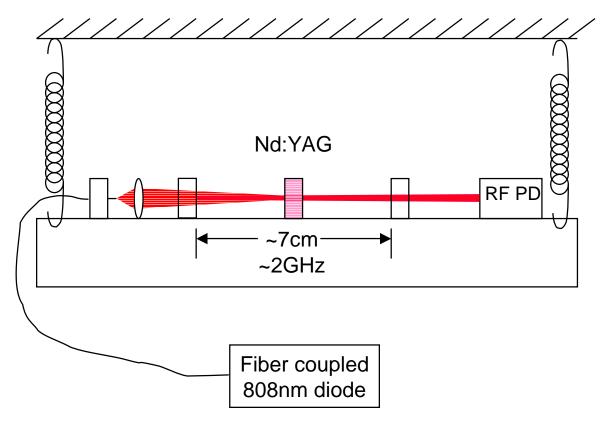


Spectrum analyzer trace of the RF beat frequency ~500 Hz linewidth at 5GHz



Proposed two-mode laser

 Laser cavity, pump source, and RF photodiode are all mounted on a single block and seismically isolated inside a vacuum chamber.





Proposed Work

Year 1 milestones

- Contruct CO₂ and Nd:YAG two-frequency lasers
- Identify technical noise sources (and eliminate if possible)
- Measure Allan variance of laser RF sources
- Optimize laser operation for minimum Allan variance

Year 2 milestones

- Construct second set of laser sources with upgrades
- Measure absolute phase noise spectra
- Continued optimization and identification of noise sources